

Uitwerking Toets QM 1

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- 1)
 a. $\hat{H}(x, t)\Psi(x, t) = i\hbar \frac{\partial \Psi(x, t)}{\partial t}$
 b. $V \neq V(t)$
 c. $\Psi(x, t) \equiv \varphi(x)f(t)$

$$\hat{H}(x)\varphi(x)f(t) = i\hbar\varphi(x)\frac{f(t)}{dt} \rightarrow \frac{1}{\varphi}\hat{H}\varphi = \frac{i\hbar}{f}\frac{df}{dt} \equiv E$$

$$\rightarrow \frac{df}{dt} = \frac{E}{i\hbar}f \rightarrow f \exp\left(\frac{-iEt}{\hbar}\right) \quad \text{met } E \text{ oplossing van } \hat{H}\varphi = E\varphi$$

- 2) a. $[\hat{H}, \hat{Q}] = 0 \rightarrow \frac{d}{dt}\langle Q \rangle = 0 \rightarrow \langle Q \rangle$ behouden.
 b.

$$\left. \begin{aligned} \frac{d}{dt}\langle p_x \rangle &= \frac{i}{\hbar}\langle [\hat{H}, \hat{p}_x] \rangle \\ [\hat{H}, \hat{p}_x] &= \left[\frac{\hat{p}_x^2}{2m} + \hat{V}, \hat{p}_x\right] = [\hat{V}, \hat{p}_x] = -\frac{\hbar}{i}\left(\frac{dV}{dx}\right) \end{aligned} \right\} \Rightarrow$$

$$[\hat{V}, \hat{p}_x]\psi = V\frac{\hbar}{i}\frac{d\psi}{dx} - \frac{\hbar}{i}(V\psi) = -\frac{\hbar}{i}\left(\frac{dV}{dx}\right)\psi$$

$$\Rightarrow \frac{d}{dt}\langle p_x \rangle = \left\langle -\left(\frac{dV}{dx}\right) \right\rangle$$

c. Verwachtingswaarden beschouwen als klassieke observabelen geeft $\frac{d}{dt}p_x = -\frac{dV}{dx} = F_x$ d.w.z. 2^e wet van Newton.

- 3) a. $\hat{H} = -\frac{\hbar^2}{2m}\frac{d^2}{dx^2}$

$$\hat{H}\psi_n(x) = -\frac{\hbar^2}{2m}\frac{d^2}{dx^2}\left(\sqrt{\frac{2}{a}}\sin\frac{n\pi x}{a}\right) = \frac{\hbar^2}{2m}\frac{n^2\pi^2}{a^2}\sqrt{\frac{2}{a}}\sin\frac{n\pi x}{a}$$

$$= \frac{n^2\pi^2\hbar^2}{2ma^2}\psi_n(x) \rightarrow \text{eigenwaarden } n^2\frac{\pi^2\hbar^2}{2ma^2} \quad (n = 1, 2, \dots)$$

b. Deeltje in ψ_1 : $\psi_1(x) = \sqrt{\frac{2}{a}}\sin\frac{\pi x}{a}$ voor $0 \leq x \leq a$
 $\psi_1(x) = 0$ voor $a < x \leq 2a$

Eigentoestanden \hat{H} na verplaatsen $\equiv \psi'_n(x) = \sqrt{\frac{1}{a}}\sin\frac{n\pi x}{2a}$ voor $0 \leq x \leq 2a$.

Iedere toestand, dus ook ψ_1 te schrijven als lin. comb. van ψ'_n ($n = 1, 2, \dots$) \longrightarrow
 $\psi_1 = \sum_n c_n \psi'_n \longrightarrow$ kans om deeltje in ψ'_1 aan te treffen $= |c_1|^2 = |\langle \psi'_1 | \psi_1 \rangle|^2$.

$$\langle \psi'_1 | \psi \rangle = \int_0^{2a} \psi'_1 \psi_1 dx = \int_0^a \sqrt{\frac{1}{a}} \sqrt{\frac{2}{a}} \sin \frac{\pi x}{2a} \sin \frac{\pi x}{a} dx = \frac{4\sqrt{2}}{3\pi}$$

$$|c_1|^2 = \frac{32}{9\pi^2} (\approx 0,36).$$

4) a.

$$\left. \begin{aligned} \hat{a}_+ \hat{a}_- \psi_n &= \hat{a}_+ \sqrt{n\hbar\omega} \psi_{n-1} = n\hbar\omega \psi_n \\ \hat{a}_- \hat{a}_+ \psi_n &= \hat{a}_+ \sqrt{(n+1)\hbar\omega} \psi_{n+1} = (n+1)\hbar\omega \psi_n \end{aligned} \right\} [\hat{a}_+, \hat{a}_-] = -\hbar\omega$$

b.

$$\frac{1}{2}\hbar\omega \begin{pmatrix} 1 & & & \\ & 3 & & \\ & & 5 & \\ & & & \ddots \end{pmatrix} \text{matrixrepresentatie is diagonaal in de basis van eigenvectoren.}$$

$$\text{c. } \Psi(x, t) = \sum_{n=0}^{\infty} c_n \psi_n(x) \exp\left(\frac{-iE_n t}{\hbar}\right).$$

d.

$$\langle \Psi(x, t) | \hat{Q} \Psi(x, t) \rangle = \sum_{n,m=0}^{\infty} c_n^* c_m \langle \psi_n | \hat{Q} \psi_m \rangle \exp\left(\frac{-i(E_m - E_n)t}{\hbar}\right) =$$

$$\{E_n - E_m = (n - m)\hbar\omega \equiv p\hbar\omega; \quad n, m \in [0, \infty] \rightarrow p \in [-\infty, \infty]\}$$

$$= \sum_{p=-\infty}^{+\infty} A_p \exp(ip\omega t) \quad \text{met} \quad A_p = \sum_{n=0}^{\infty} c_n^* c_{n-p} \langle \psi_n | \hat{Q} \psi_{n-p} \rangle$$

$$\text{e. } \hat{a}_+ - \hat{a}_- = \frac{1}{\sqrt{2m}}(2im\omega\hat{x})$$

$$\langle \psi_n | \hat{x} | \psi_m \rangle = \frac{1}{i\sqrt{2m\omega}} \langle \psi_n | (\hat{a}_+ - \hat{a}_-) \psi_m \rangle = \frac{1}{i\sqrt{2m\omega}} \left(\sqrt{(m+1)\hbar\omega} \underbrace{\langle \psi_n | \psi_{m+1} \rangle}_{=\delta_{n,m+1}} - \sqrt{m\hbar\omega} \underbrace{\langle \psi_n | \psi_{m-1} \rangle}_{=\delta_{n,m-1}} \right)$$

$$A_p = \frac{\sqrt{\hbar}}{i\sqrt{2m\omega}} \sum_{n=0}^{\infty} c_n^* c_{n-p} \left(\sqrt{n-p+1} \delta_{n,n-p+1} - \sqrt{n-p} \delta_{n,n-p-1} \right)$$

$$A_1 = \frac{\sqrt{\hbar}}{i\sqrt{2m\omega}} \sum_{n=0}^{\infty} (c_n^* c_{n-1} \sqrt{n})$$

$$A_{-1} = \frac{\sqrt{\hbar}}{i\sqrt{2m\omega}} \sum_{n=0}^{\infty} (-c_n^* c_{n+1} \sqrt{n+1})$$

$$A_p = 0 \forall p \neq \pm 1$$